CHEMISTRY OF MATERIALS

Luminescence and Microstructural Features of Eu-Activated LiBaPO₄ Phosphor

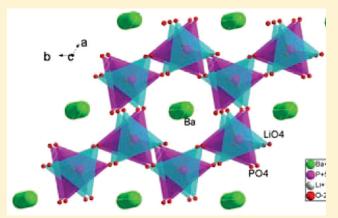
Suyin Zhang,⁺ Yosuke Nakai,[‡] Taiju Tsuboi,[‡] Yanlin Huang,^{*,†} and Hyo Jin Seo^{*,§}

⁺College of Chemistry, Chemical Engineering and Materials Science, Soochow University, Suzhou 215123, China

[‡]Faculty of Engineering, Kyoto Sangyo University, Kamigamo, Kyoto 603-8555, Japan

[§]Department of Physics, Pukyong National University, Busan 608-737, Republic of Korea

ABSTRACT: Eu²⁺-activated LiBaPO₄ phosphor was synthesized by conventional solid-state reaction. The photoluminescence excitation and emission spectra, the temperature dependent luminescence intensities (12–450 K), and decay curves of the phosphor were investigated. With the increasing of temperatures, the emission bands of LiBaPO₄:Eu²⁺ show the abnormal blue-shift and the decreasing of emission intensity. The natures of the Eu²⁺ emission in LiBaPO₄, for example, the luminescence quenching temperature, and the activation energy for thermal quenching (ΔE), were reported. The afterglow fluorescence was detected in LiBaPO₄:Eu²⁺ phosphor. Together with the Eu²⁺ luminescence, Eu³⁺ ions with the abnormal crystal field were observed. The site-selective excitation in the ⁵D₀ \rightarrow ⁷F₀ region for Eu³⁺ ions, emission spectra, and decay curves have been investigated using a pulsed, tunable, and narrowband dye laser to



detect the microstructure and crystallographic surrounding of $Eu^{3+,2+}$ at Ba^{2+} sites in LiBaPO₄. The multiple sites structure of Eu^{2+} and Eu^{3+} ions in LiBaPO₄ lattices was suggested. The lower quenching temperature, afterglow, and luminescence mechanism were discussed. The photoluminescence quantum efficiencies of LiBaPO₄: Eu^{2+} were measured and compared with the reported phosphors. Different from the published data on LiBaPO₄: Eu^{2+} , this investigation indicates that LiBaPO₄: Eu^{2+} is not a good phosphor candidate applied in white light emitting diode.

KEYWORDS: luminescence, optical materials and properties, phosphors, LED, tridymite structure

INTRODUCTION

The phosphates with ABPO₄ formula (A and B are mono- and divalent cations, respectively) are in a large family of monophosphates with the different structure types strictly depending on the relative size of the A and B ions.^{1,2} These compounds have been considered to be efficient luminescent hosts due to their excellent thermal and hydrolytic stability.^{3,4} Recently, Eu²⁺-doped ABPO₄ phosphates have received much attention for the potential applications as new white light emission diodes (W-LEDs) phosphors,⁴ such as KCaPO₄:Eu^{2+,5} KSrPO₄:Eu^{2+,6} KBaPO₄:Eu^{2+,7,8} LiSrPO₄:Eu^{2+,9} NaCaPO₄:Eu^{2+,10} and ABaPO₄:Eu²⁺ (A = Na, K).¹¹

ABPO₄ has a tridymite structure (β -SiO₂) when B is rather small and A is large,² for example, LiBaPO₄. This structure involves the existence of a 1:1 ordering between the cross-linked LiO₄ and PO₄ tetrahedra on the basis of an isotypy with the β -SiO₂ trydimite.¹² Eu²⁺-doped compounds with tridymite structure have been confirmed to be notable for their good luminescent performance, for example, MAl₂O₄ (M = Ba, Sr, and Ca)^{13,14} and BaMgSiO₄.^{15,16} There are two barium sites Ba(1) and Ba(2) with nine-coordination in BaAl₂O₄ lattices. Therefore, two Eu²⁺ emission bands were observed in Eu²⁺-doped BaAl₂O₄.^{14,17} In the structure of BaMgSiO₄ there are three different Ba sites in equal amounts in the lattice. Two sites Ba(1) and Ba(2) are coordinated by nine oxygen ions, whereas the third Ba(3) is surrounded by six oxygen ions.¹⁸ Two Eu(1) and Eu(2) in the structure of BaMgSiO₄ are comparable with those in BaAl₂O₄,¹⁷ which give emissions at long wavelength, while the Eu(3) site shows a short emission band at 398 nm.¹⁵ In BaAl₂O₄ and BaMgSiO₄, the Eu²⁺ luminescence is at low energies due to preferential orientation of a d orbital of Eu²⁺, no preferential orientation of the d orbital is left due to the strongly distorted stuffed-tridymite lattice. Consequently, only a blue emission at 440 nm has been observed.^{13,19} This indicates that although these compounds have the similar stuffed-tridymite-type structure, the chemically different compounds have the important influence on the Eu²⁺ luminescence.¹⁵

It has been confirmed that $LiBaPO_4$ is suitable for second harmonic generating (SHG) responses and shows larger SHG

Received:October 4, 2010Published:January 19, 2011

effects relative to quartz.²⁰ The luminescence properties of Ce³⁺, Sn²⁺, and Cu⁺ doped LiBaPO₄ have been reported.^{12,21,22} Waite²³ first report the Eu²⁺ luminescence doped in LiBaPO₄. Just as in BaAl₂O₄, Eu²⁺ ions on barium sites in LiBaPO₄ have a preferential orientation of the d orbital. Recently, Wu et al. have reported that LiBaPO₄:Eu²⁺ is a good candidate phosphor applied in W-LEDs and exhibits higher thermal stability than that of commercially available Y₃Al₅O₁₂:Ce³⁺.²⁴

In this paper, we present an insightful investigation of the luminescence of LiBaPO₄:Eu²⁺. We obtained the different results from those reported in the references.²⁴ In addition, the luminescence quantum efficiency, crystallographic sites, and the microstructure of Eu²⁺ ions doped in LiBaPO₄ are reported. The temperature dependent luminescence intensities and decay curves (12-450 K) were measured. The activation energy for thermal quenching was obtained for LiBaPO₄:Eu²⁺. LiBaPO₄:Eu²⁺ has a lower quenching temperature, afterglow fluorescence, and changing color with increase of temperature. The characteristics indicate that this phosphor is not a candidate applied in W-LEDs. The excitation spectra have been investigated in the ${}^{5}D_{0} \rightarrow {}^{7}F_{0}$ region for Eu³⁺ ions in LiBaPO₄ by using a pulsed, tunable, and narrowband dye laser to detect the Ba^{2+} cation sites in LiBaPO₄. The multiple sites structure of Eu²⁺ ions in LiBaPO₄ lattices was suggested and discussed.

EXPERIMENTAL SECTION

Polycrystalline samples of Eu²⁺-doped LiBaPO₄ were synthesized using a conventional solid-state reaction. The starting material was a stoichiometric mixture of reagent grade BaCO₃, Li₂CO₃, (NH₄)₂HPO₄ (A.R. grade), and Eu₂O₃ (99.99% purity). First, the mixture was heated up to 350 °C in 10 h and kept at this temperature for 5 h. The obtained powder was mixed and then heated up to 750 °C for 5 h in air. After that, the sample was thoroughly mixed and heated for 10 h in a reduction atmosphere (N₂:H₂ = 95:5) at 1050 °C obtained by DTA analysis. The Eu²⁺ doped samples were prepared with the doping concentration of 5% of BaCO₃. The XRD pattern was collected on a Rigaku D/Max diffractometer operating at 40 kV and 30 mA with Bragg–Brentano geometry using Cu Kα radiation (λ = 1.5405 Å).

The photoluminescence excitation and emission spectra were recorded on a Perkin-Elmer LS-50B luminescence spectrometer with Monk-Gillieson type monochromators and a xenon discharge lamp used as excitation source. Luminescence spectra were recorded between 10 and 300 K on a 75 cm monochromator (Acton Research Corp. Pro-750) equipped with a helium flow cryostat and observed with a photomultiplier tube (PMT) (Hamamatsu R928). To study the thermal quenching from 20 to 500 °C, the same spectrofluorimeter was equipped with a homemade heating cell under the excitation of a 365 nm UV lamp. The luminescence decay was measured using the third harmonic (355 nm) of a pulsed Nd:YAG laser. The crystallographic site of Eu³⁻ ions in this host was investigated by the site-selective excitation measurement (for the description of the experimental method see ref 25). Quantum efficiency (QE) was measured by an Absolute Photoluminescence Quantum Yield Measurement System (C9920-02, Hamamatsu) at room temperature. The excitation was done by changing the excitation wavelength of light from a 150 W Xe lamp.

RESULTS AND DISCUSSION

Phase Formation. The X-ray diffraction (XRD) patterns of LiBaPO₄: Eu^{2+} together with the Joint Committee on Powder Diffraction Standards (JCPDS) card No. 14-0270 are shown in Figure 1. From a comparison between them, the position and

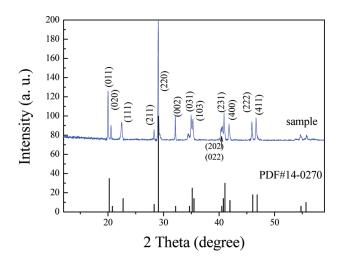


Figure 1. XRD pattern of LiBaPO₄: Eu^{2+} of this work and JCPDs card No.14-0270.

intensity of the peaks are the same. No impurity lines were observed, and the entire pattern could be well indexed to a $LiBaPO_4$ single phase.

The material LiBaPO₄ belongs to the stuffed tridymite structure with a hexagonal unit cell and space group $P6_3$.²⁶ In this stuffed tridymite structure, each PO₄ tetrahedron is connected with four LiO₄ tetrahedrons and each LiO₄ tetrahedron is connected with four PO₄ tetrahedrons, and the two types of tetrahedrons are linked by corner sharing and build up the tridymite skeleton. As shown in Figure 2, the three-dimensional network of corner-sharing LiO₄ and PO₄ tetrahedrons forms channels in the *a*- and *c*-directions, where Ba²⁺ ions are located (Figure 2 a,b). In this structure, there are three types of Ba²⁺ sites, Ba(1), Ba(2), and Ba(3). The structure of LiBaPO₄ is closely related to BaAl₂O₄. The most important difference between them is the size of the two kinds of tetrahedra in LiBaPO₄ ([LiO₄]⁷⁻ and [PO₄]³⁻), whereas there is only one kind of tetrahedron in BaAl₂O₄ ([A1O₄]⁵⁻).

Photoluminescence Spectra. The excitation spectra of Li-BaPO₄:Eu²⁺ in Figure 3 by monitoring 473 and 530 nm show a broad peak from 230 to 450 nm, which can be attributed to 4f-5d transition of Eu²⁺ ions. It suggests that LiBaPO₄:Eu²⁺ could be effectively excited by UV chips (360-400 nm). The emission spectra of the Eu²⁺-doped LiBaPO₄ under the excitation of 280 and 385 nm in Figure 3 show one emission band peaking at 485 nm, which is ascribed to the 4f⁶5d \rightarrow 4f⁷ (⁸S_{7/2}) transition on Eu²⁺ ions. The phosphor presents bright green luminescence. These results agree with what was reported by Wu et al.^{24,27}

The asymmetric emission spectra (λ_{ex} = 280 and 385 nm) in Figure 3 show that Eu²⁺ has more than one emission center in LiBaPO₄ lattices, which can be deconvoluted into at least two Gaussian components peaked at 473 and 530 nm. The emission spectrum of LiBaPO₄:Eu²⁺ shows no dependence of the emission spectra on the excitation wavelengths. However, the excitation spectra for the two emission bands have a noticeable difference. The rough estimation of the Stokes shifts for the emissions at 473 and 530 nm are about 5400 cm⁻¹, and 9700 cm⁻¹, respectively.

In the stuffed-tridymite-type structures, Ba ions have Ba(1)and Ba(2) in nine-coordination in $BaAl_2O_4$ or three sites Ba(1)and Ba(2) in nine-coordination and Ba(3) in six-coordination in

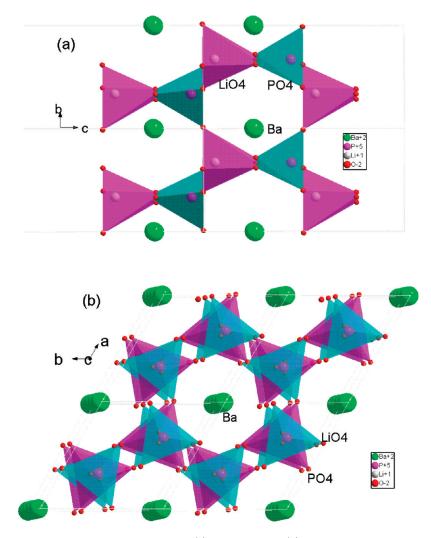


Figure 2. Schematic views of the structure of LiBaPO₄ along the *a*- (a) and *c*-directions (b).

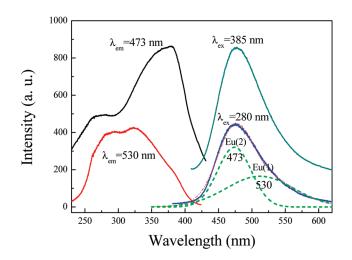


Figure 3. Excitation ($\lambda_{em} = 473$ and 530 nm) and emission spectra ($\lambda_{ex} = 280$ and 385 nm) of LiBaPO₄:Eu²⁺. One emission spectrum is given by a Gaussian fit into two components.

BaMgSiO₄. The so-called Ba(1) and Ba(2) are distinguished by the average distance from the oxygen ions, for example, Ba(1) $-O^{2^-}$ is 2.89 Å and Ba(2) $-O^{2^-}$ is 2.94 Å in BaMgSiO₄.¹⁸ Usually, Eu(1) and Eu(2) give emissions at long wavelength, and a short emission band is from Ba(3). For example, in BaMgSiO₄ there are three Eu²⁺ emissions at 502, 509, and 398 nm.¹⁶ When the crystal environments are analogous, the Eu²⁺ center with a shorter Eu²⁺ $-O^{2-}$ distance will give a longer wavelength emission. Since the crystal environments of Eu²⁺ ions on Ba(1) and Ba(2) sites are very similar in BaMgSiO₄, the 502 nm emission was ascribed to the Ba(2) site, and the 509 nm to Eu²⁺ ions on the Ba(1) sites. The much higher energy of Eu²⁺ emission at 398 nm on Ba(3) sites could be explained by means of preferential orientation of a d orbital of the Eu²⁺ ion.

The luminescence spectrum at 12 K (Figure 4) consists of a broad band at 480 nm and a band at about 530 nm; in addition a small shoulder at 412 nm can be observed. At temperatures above 297 K the two longer wavelength bands are only visible, and the 412 nm emission band disappears. Analogously, it is reasonable that the Eu^{2+} emission at 473 and 530 nm could be ascribed to Ba(2)(Eu2) and Ba(1)(Eu1), respectively. And 412 nm emission is from Ba(3)(Eu3) sites, which quenches at RT similar to that in $BaMgSiO_4$. These three occupations will be confirmed by site-selective luminescence spectra of Eu^{3+} mentioned below.

Photoluminescence Quantum Efficiency (QE). The QE of a phosphor is important for its application. The QE of $LiBaPO_4$:

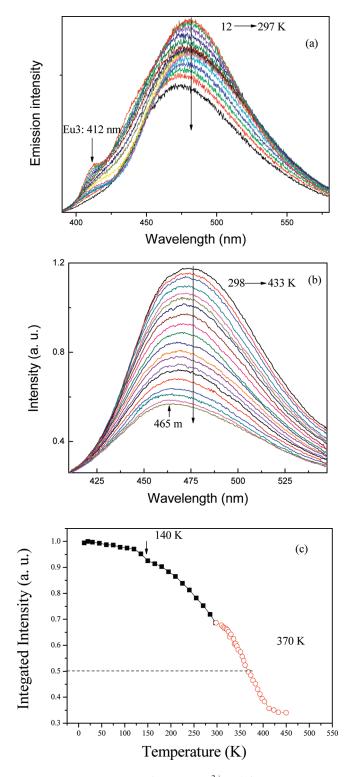


Figure 4. Emission spectra of LiBaPO₄: Eu^{2+} at different temperatures (a and b) and the temperature dependence of the integrated emission intensity normalized with respect to the value at 12 K (c).

 ${\rm Eu}^{2+}$ luminescence was measured to be 58.5% at the excitation of 330 nm light at 300 K. However, LiBaPO₄:Eu²⁺ has a lower luminescence QE value in comparison with the reported results in some LED phosphors in Table 1. Certainly, the QE usually depends on the synthesis conditions and Eu²⁺ concentrations. The higher quantum yields can be obtained by further improving

Table 1. Quantum Efficiencies at Room Temperature of	
LiBaPO ₄ : Eu^{2+} and the Reference Phosphors	

phosphors	λ_{ex}	$\lambda_{ m em}$	QE (%)	
LiBaPO ₄ :Eu ²⁺	330 nm	483 nm	58.5	this work
$CaSrS:Eu^{2+}$ (1 mol %)	350 nm	663 nm	53.4	ref 28
AlN:Eu ²⁺ (0.10 mol %)	365 nm	465 nm	63.0	ref 29
$BaMgAl_{10}O_{17}:Eu^{2+} (2 mol \%)$	254 nm	450 nm	70.0	ref 30
YAG:Ce ³⁺ (mol 3%)	460 nm	560 nm	70.0	ref 31
$SrSi_2N_2O_2:Eu^{2+} (2 \text{ mol } \%)$	450 nm	535 nm	91.0	ref 32

the synthesis conditions to reduce the number of defects and impurities.

Temperature Dependent Luminescence and Thermal Quenching. In general, the temperature dependence of W-LED phosphors is important because it has great influence on the light output and color rendering index. Phosphors must sustain emission efficiency at temperatures of about 150 °C over a long term when they are used in white LEDs. It is thus required that the thermal quenching of phosphors should be small, typically for high-power ones. The thermal quenching of LiBaPO₄:Eu²⁺ was evaluated by measuring the temperature-dependent emission intensities as shown in Figure 4.

The emission intensities decrease as the temperature increases as shown in Figure 4a,b. Figure 4c represents the temperature dependence of the integrated emission intensity normalized to 12 K. The thermal quenching begins during the heating at 140 K. The thermal quenching temperature, $T_{0.5}$, defined as the temperature at which the emission intensity is 50% of its original value, is 370 K. It can be seen from references that the $T_{0.5}$ for Eu²⁺-doped BaSi₂O₅ is 460 K,³³ that for Cs₂MP₂O₇ (M = Ca, Sr) is 600 K,³³ that for Sr₃(PO₄)₂ and Ba₃(PO₄)₂ is >550 K,³⁴ that for SrSi₂O₂N₂ is 600 K,³⁵ and that for KBPO₄:Eu²⁺ (B = Sr, Ba) is 673 K.⁴ The $T_{0.5}$ for iso-structural BaMgSiO₄: Eu²⁺ is 300 K.¹⁵ In comparison with the reported phosphors, LiBaPO₄:Eu²⁺ has a relatively lower quenching temperature.

P. Dorenbos has proposed that the main mechanism responsible for the thermal quenching of Eu^{2+} luminescence in solids is the ionization of the electron from the lowest energy level of the relaxed Eu^{2+} 4f⁶5d¹ electronic configuration to the host lattice conduction band level.³⁶ Following this suggestion, it is responsible that the thermally activated ionization from the 4f⁶5d state is for temperature quenching of the luminescence in LiBaPO₄: Eu^{2+} . The temperature dependence of the luminescence intensity is described by a modified Arrhenius equation as following equation:³⁷

$$I_T = \frac{I_0}{1 + c \exp\left(-\frac{\Delta E}{kT}\right)} \tag{1}$$

where I_0 is the initial emission intensity, I_T is the intensity at different temperatures, ΔE is activation energy of thermal quenching, *c* is a constant for a certain host, and *k* is the Boltzmann constant (8.629 × 10⁻⁵ eV). The activation energy is the energy required to raise the electron from the relaxed excited level into the host lattice conduction band. Figure 5 plots $\ln[(I_0/I_T) - 1]$ versus 1000/*T* for LiBa_{0.95}Eu_{0.05}PO₄. According to eq 1, the activation energy ΔE was calculated to be 0.03 eV. Moreover, in Figure 4a,b the wavelength of the maximum emission shifts toward blue region from 480 to 465 nm with increasing temperature. This

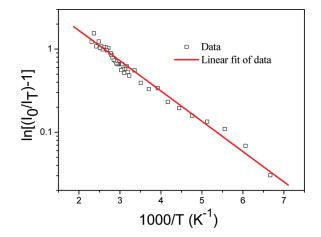


Figure 5. Activation energy of the thermal quenching of LiBaPO₄:Eu²⁺ phosphor fitted in eq 1.

shift with increasing temperature is indicative of the chromatic instability of $LiBaPO_4:Eu^{2+}$.

Temperature Dependent Luminescence Lifetimes. Figure 6a shows the luminescence decay curves of LiBaPO₄:Eu²⁺ at different temperatures. The decays display nearly exponential curves. The luminescence lifetimes of the Eu²⁺ ions are calculated as a function of temperature and displayed in Figure 6b. On the base of the suggested mechanism for the thermal quenching of Eu²⁺ is the ionization of the electron from the lowest energy level of the relaxed Eu²⁺ 4f⁶5d¹ electronic configuration to the host lattice conduction band level.³⁶ The activation energy (ΔE) for thermal quenching of the Eu²⁺ was also determined by measuring the temperature dependence of the Eu²⁺ emission lifetime. The drawn line in the figure is fitted to the equation as³³

$$\tau(T) = \frac{\tau_{\rm r}}{1 + [\tau_{\rm r}/\tau_{\rm nr}] \exp(-\Delta E/kT)}$$
(2)

where k is the Boltzmann constant and $\tau_{\rm r}$ and $\tau_{\rm nr}$ are radiative and nonradiative decay times. The thermal activation energy for thermal quenching (ΔE) is fitted to be 0.032 eV. This is matched with the calculation in eq 1.

In Figure 6b, the luminescence lifetime of Eu^{2+} is near a constant (around 820 ns) from 12 K up to 240 K and then drops at higher temperature, presenting a typical temperature quenching behavior. The rapid shortening of the Eu^{2+} emission lifetime for T > 240 K indicates the onset of nonradiative transitions, which may be a thermal ionization and/or energy transfer process. In the case of thermal ionization, the electron in the excited 5d state of the Eu^{2+} auto-ionizes into the conduction band and is delocalized. This is responsible for the observed quenching. The calculated activation energy of 0.03 eV indicates that the thermal energy required to promote an electron from the lowest excited state of Eu^{2+} to the host lattice conduction band is not high.³⁶

As shown above, LiBaPO₄:Eu²⁺ does not show the stable luminescence at high temperature. This is not in agreement with the reported results by Wu et al.:^{24,27} LiBaPO₄:Eu²⁺ with the $T_{0.5}$ of 450 K has higher thermal stability than commercially available YAG:Ce³⁺. However, our results are similar to that recently reported for LiSrPO₄:Eu²⁺ phosphor by Lin et al.:⁴ LiSrPO₄:Eu²⁺ has worse thermal stability than KBaPO₄:Eu²⁺ and KSrPO₄:Eu²⁺. This was explained that Li⁺ promotes the

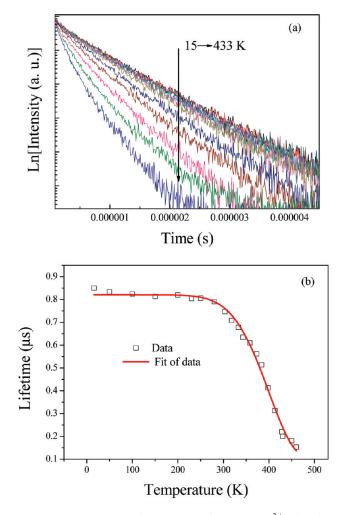


Figure 6. Luminescence decay curves of LiBaPO₄:Eu²⁺ phosphor under the excitation of 355 nm of a pulsed Nd:YAG laser at 15 to 433 K (a) and the temperature dependent lifetimes of the Eu²⁺ emission fitted into eq 2 (b).

conversion of Eu^{2+} to Eu^{3+} more than K^+ ion does, based on electronic affinity.

A low $T_{0.5}$ and low activation energy of LiBaPO₄:Eu²⁺ might be related to the lowest energy level of the Eu²⁺ 4f⁶5d¹ electronic configuration not being well isolated from the host lattice conduction band. Recently, Lin et al.⁴ have elucidated the luminescent intensity of activators in ABPO₄ phosphors at various temperatures with reference to crystal structure and the coordination environment. In this case, the microstructure and crystallographic surrounding of Eu²⁺ ions at Ba²⁺ sites should be understood.

Afterglow and Emission Spectra from Eu^{3+} . The afterglow in LiBaPO₄:Eu²⁺ can be obviously seen after switching off the excitation of a UV lamp. The afterglow decay curve (Figure 7) is detected after the irradiation of UV light (254 nm). The curve can be fitted into a biexperimental decay function, indicating at least two decay processes: a fast decay (15 s) and a slow decay (79 s). It is probably due to the existence of traps with appropriate depth, and the related work is being carried on. These results suggest that electrons and holes produced by UV excitation move back to Eu^{2+} sites in the crystals through thermal hopping and tunneling and recombine radiatively at Eu^{2+} sites. Such afterglow was also reported in Eu^{2+} -doped Ba_2SiO_4 and Ba_3SiO_5 crystals.³⁸

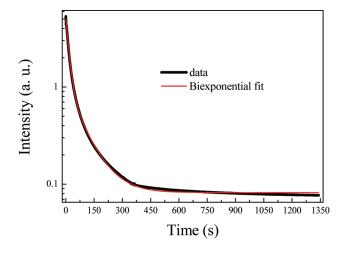


Figure 7. Experimental afterglow curve of LiBaPO₄:Eu²⁺ phosphor.

Actually, in LiBaPO₄:Eu²⁺, besides the Eu²⁺ ions, the emission from Eu³⁺ also existed. From the time-resolved spectra, the Eu³⁺ luminescence can be distinctly displayed as shown in Figure 8a. Under the time delay of 50 μ s, the emission transitions ${}^{5}D_{0}-{}^{7}F_{0,1,2}$ from Eu³⁺ ions can be observed. And the stronger Eu³⁺ emission can be observed with the decrease of Eu²⁺ emission by monitoring after a long delay time after laser excitation. This indicates that actually two different valence states, +2, and +3, are available for Eu in LiBaPO₄:Eu²⁺. Figure 8b shows the very different lifetimes for the Eu²⁺ and Eu³⁺ ions; therefore, the time-resolved spectra can separate their emission.

It should also be remarked that although the LiBaPO₄:Eu²⁺ sample was obtained in a reduction atmosphere, this reduction could not be completely realized. Under reducing conditions, Eu²⁺ ions can be stable on alkaline earth metal sites. However, when Eu²⁺ ions are heated within a certain temperature range, they may be oxidized as $Eu^{2+} \rightarrow Eu^{3+} + e^-$. Therefore, a luminescence of the Eu³⁺ can exist. The phenomenon is related to the crystal structure and the coordination environment of the activators at different temperatures.

In oxide matrixes the $(4f)^n \rightarrow (4f)^{n-1}(5d)^1$ transitions usually occur in the UV spectral region for Ln^{3+} ions and in the visible region for Ln^{2+} ions. They are parity allowed, and therefore their intensity is many times higher than that of the 4f-4f transitions that may be masked completely. Therefore, the PL spectra do not show any signature of Eu³⁺ in the sample under the excitation of continued wavelength, which means the Eu³⁺ luminescence is suppressed. Tang and Chen¹¹ recently reported this phenomenon in Eu²⁺ doped ABaPO₄:Eu (A = Na⁺, K⁺). XPS results show that Eu²⁺ and Eu³⁺ ions coexist in these compounds even though the spectra of Eu³⁺ were not observed in that work.

The emission spectra of Eu³⁺ doped in LiBaPO₄ shown in Figure 8a present the abnormal features: the unusual high energy position (573 nm, 17 452 cm⁻¹) and the abnormal relative intensity (many times stronger than the ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$) for the forbidden ${}^{5}D_{0} \rightarrow {}^{7}F_{0}$ transition. Usually, this observation is connected with the charge-compensating oxide ion and experience of an abnormal crystal field (CF) around the Eu³⁺ ions.³⁹ In such a situation an unusual crystal field is expected, affecting both the energy-level location and the relative intensity of the radiative transitions from the ${}^{5}D_{0}$ level.

Site-Selective Spectra and Discussions. BaLiPO₄ belongs to the stuffed tridymite structures which are derived from the

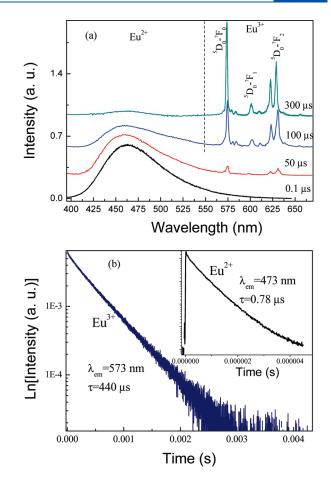


Figure 8. Time-resolved spectra LiBaPO₄: Eu^{2+} measured at different delay times after the laser excitation as labeled in part (a) and the luminescence decay curves from Eu^{2+} and Eu^{3+} in LiBaPO₄: Eu^{2+} (b).

structure of β -SiO₂. The structure consists of tetrahedral SiO₄ units which are linked together through corner sharing of all four oxygen atoms of each tetrahedron, thereby creating a one-dimensional six-ring channel. P⁵⁺ and Li⁺ replace Si⁴⁺ in the tetrahedra of the SiO₂ tridymite; Ba²⁺ will occupy sites in channels parallel to the *c*-axis.²³ In the tetrahedra network, each PO₄ shares four corners with LiO₄ tetrahedra. Thus, the tetrahedra form a six-membered ring, and Ba cations occupy the channels running through the tetrahedral network. Although the crystal structure of BaLiPO₄ is similar to trydimite, the direction of tetrahedra is different.

It is interesting to further investigate whether or not the abnormal crystal field (CF) around the Eu^{3+} ions can be kept in LiBaPO₄ prepared in air atmosphere. Eu^{3+} (4f⁶ configuration) is largely used as a probe of the surrounding symmetry because the emission and excitation lines between ${}^{5}D_{0}$ and ${}^{7}F_{0}$ levels are nondegenerate.⁴⁰ By site-selective excitation and emission spectra, the presence of crystallographic nonequivalent sites in a given host matrix can be revealed, for example, three Eu^{2+} sites in LiBaPO₄ lattices. This is helpful to elucidate the luminescence mechanism of the Eu^{2+} ion by revealing the microstructure and crystallographic surrounding.

The Eu sites in LiBaPO₄ were revealed by the excitation spectra of the ${}^{7}F_{0} \rightarrow {}^{5}D_{0}$ transition of Eu³⁺ ions in LiBaPO₄ obtained by heating the same sample of LiBaPO₄:Eu²⁺ in air atmosphere. Figure 9 shows the excitation spectra corresponding

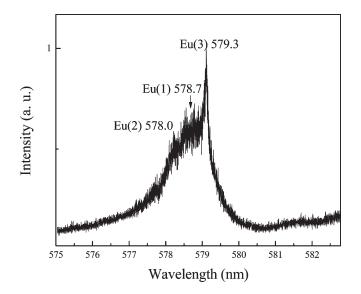


Figure 9. Excitation spectra for ${}^{7}F_{0} \rightarrow {}^{5}D_{0}$ transition of Eu³⁺ in LiBaPO₄:Eu³⁺ by monitoring the total luminescence.

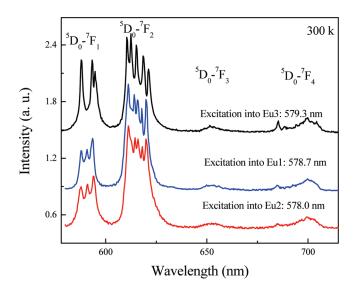


Figure 10. Site-selective emission spectra of ${}^{5}D_{0} \rightarrow {}^{7}F_{J}$ (J = 1, 2, 3, 4) for LiBaPO₄:Eu³⁺ by exciting into site Eu(1): 578.0 nm, Eu(2): 578.7 nm, and Eu(3): 579.3 nm.

to the $^7F_0 \rightarrow {}^5D_0$ transition, obtained by monitoring the total luminescence of Eu $^{3+}$ ions. The $^7F_0 \rightarrow {}^5D_0$ excitation spectra consist of the transition lines at 578.0, 578.7, and 579.3 nm. Usually in an ionic system, the position of the ${}^5D_0 \rightarrow {}^7F_0$ transition level of Eu $^{3+}$ depends on the ionicity of the system and the size and coordination of the Eu $^{3+}$ or Sm $^{2+}$ site. The high coordination and the long metal—oxygen atomic distance for Eu $^{3+}$ in the lattice are responsible for the high energy of the 5D_0 level. 41

In the stuffed-tridymite-type structures, the crystal environments of Eu²⁺ ions on Ba1 and Ba2 sites are similar and the Ba(2) $-O^{2-}$ distance is longer than that of Ba(1) $-O^{2-}$.¹⁸ Therefore, the $^7F_0 \rightarrow ^5D_0$ transitions at 578.0 and 578.7 are ascribed to Ba(2)(Eu2) and Ba(1)(Eu1), respectively. And Ba(3)(Eu3) sites are attributed to the 579.3 nm transition line. The results in Figure 9 clearly indicate that the Eu³⁺ ions occupy three intrinsic

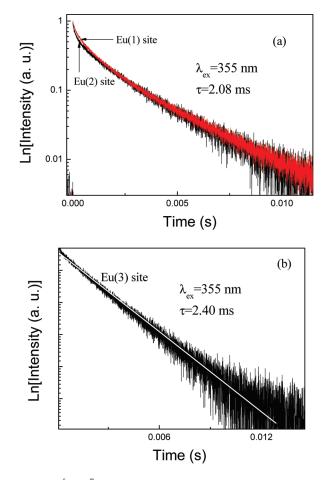


Figure 11. ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$ luminescence decay curves under the excitation of sites Eu(2): 578.0 nm and Eu(1): 578.7 nm (a) and Eu(3): 579.3 nm (b).

crystallographic sites in LiBaPO₄. Although the microstructure around the Eu^{2+} and Eu^{3+} ions doped in LiBaPO₄ could be different because of the different charge compensations in the same host, however, the conclusion of the at least three Eu^{2+} sites occupying on Ba²⁺ in LiBaPO₄ could be reasonable.

The site-selective emission spectra were recorded by tuning the laser to resonance with each excitation peak of the sites Eu(2) (578.0 nm), Eu(1) (578.7 nm), and Eu(3) (579.3 nm) (in Figure 10). The two spectra for site-selective excitation into sites Eu(2) and Eu(1) have nearly the same profile with the same emission wavelength position. The strongest emission line is the ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$ transition at 614 nm for two sites. However, the spectrum for Eu3 has a different profile. Actually, this characteristic can also be supported by the fluorescence decay curves for each site. Figure 11 shows decay curves of ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$ emission under excitation in sites Eu(2) (578.0 nm), Eu(1) (578.7 nm), and Eu(3) (580.3 nm). The decay curves for sites Eu(2) and Eu(1) exhibit the nearly same lifetime of 2.08 ms with a small difference in the initial part. The luminescence curve for Eu(3) has a good exponential decay with a longer lifetime of 2.4 ms.

It can be suggested that the Eu(1) and Eu(2) positions of the Eu^{3+} ions in LiBaPO₄ were arranged in the similar environment. This property can be explained by the high disordering of the cation positions on Eu(1) and Eu(2). This can be seen from the excitation spectra in Figure 9 and emission spectra in Figure 10; the broadening of the fluorescence lines observed can be

explained in terms of the channel in this structure which is occupied by Eu ions resulting in a larger local distortion and greater disorder over the whole structure. It seems that the Eu atoms occupy a family of sites with identical oxygen environments.

In LiBaPO₄, two Eu²⁺ sites can be detected at room temperature; however, three Eu³⁺ emissions can be observed in LiBa-PO₄:Eu³⁺. This indicates that the Eu³⁺ ions in LiBaPO₄ are more stable than the Eu²⁺ ions in the same host. The instability luminescence of Eu²⁺ ions can be also reflected by the recent reports in rare earth ion doped ABPO₄ phosphors.⁴ It has been reported that the LiSrPO₄:Eu²⁺ phosphor has worse thermal stability than KSrPO₄:Eu²⁺ and KBaPO₄:Eu²⁺. This was explained as because Li⁺ promotes the conversion of Eu²⁺ to Eu³⁺ more than the K⁺ ion. It can be suggested that in Eu²⁺ doped LiBaPO₄, Li⁺ promotes the conversion of Eu²⁺ to Eu³⁺ and the activation energy of the conversion of Eu²⁺ to Eu³⁺ is not high.

Lin et al.⁴ also reported that the LiSrPO₄:Tb³⁺ phosphors have higher luminescence stabilities than those of KSrPO₄:Tb³⁺ and KBaPO₄:Tb³⁺. They explain that trivalent Tb³⁺ in ABPO₄ forms a competitive hole-trapping center as has been confirmed by the detection of lattice defect levels. When the trivalent cation $(Tb^{3+} \text{ or } Eu^{3+})$ was doped in the divalent cation (B site) in ABPO₄, an electron is required to the charge compensation. Then LiSrPO₄:Tb³⁺ must have the best thermal stability because the Li⁺ ion captures more electrons than the K⁺ ion or Na⁺. This can explain the reason why the emissions for three Eu^{3+} sites in LiBaPO₄ are all stable at RT. As Lin et al.⁴ suggested, the thermal stabilities of all of ABPO₄ compounds could be determined from the crystal structure and the coordination environment of the rare-earth metal. Although the real quenching mechanism for Eu²⁺ doped LiBaPO₄ phosphor is not yet understood in detail, the knowledge of the site distribution over the different sites of the structure in this work would be useful to deeply analyze the luminescence quenching mechanisms in ABPO₄. To further the study, some new investigations are needed, for example, band structures and density of states (DOS) of pure LiBaPO₄, orbital distributions of atoms, the detailed coordination environment of Eu, charge transfer band, and exchange interaction acting on the Eu.

CONCLUSIONS

In conclusion, the green phosphor Eu²⁺-doped LiBaPO₄ was synthesized by the high temperature solid state method. The photoluminescence excitation and emission spectra, the temperature dependent luminescence intensities (12-450 K), and the decay curves of the phosphor were measured. The excitation of LiBaPO₄:Eu²⁺ with a broad band from 230 to 450 nm, could be effectively excited by UV chips (360-400 nm). The luminescence spectra show that Eu^{2+} ions have three emission centers in the LiBaPO₄ lattices, which are Eu(1) at 530 nm, Eu(2) at 473 nm, and Eu(3) at 415 nm. The quantum efficiency of LiBaPO₄:Eu²⁺ was 58.5% at the excitation of 330 nm. The luminescence of LiBaPO₄:Eu²⁺ shows a typical thermal quenching effect with a $T_{0.5}$ value of 370 K. From the dependences of luminescence intensities and lifetimes on temperature, the thermal activation energy of LiBaPO₄:Eu²⁺ was calculated to be 0.03 eV. Some emission features of LiBaPO₄: Eu^{2+} , that is, the afterglow, the Stokes shift, and the color shift for thermal quench ing were discussed. Three cation sites for the Eu^{2+} and Eu^{3+} emission in LiBaPO₄ host lattices were suggested. This structural feature was confirmed by the emission spectra and the excitation spectra of the ${}^7F_0 \rightarrow {}^5D_0$ transition of probe Eu $^{3+}$ ions in LiBaPO₄. The microstructure of Eu doped in LiBaPO₄ was discussed on the basis of the crystal structure and the luminescence spectra. This could be helpful for understanding the mechanisms responsible for the quenching of luminescence at high temperature and for developing new materials that have potential application for white LEDs.

AUTHOR INFORMATION

Corresponding Author

*E-mail: hjseo@pknu.ac.kr (H.J.S.), huang@suda.edu.cn (Y.H.).

ACKNOWLEDGMENT

This work was supported by Program for Postgraduates Research Innovation in University of Jiangsu Province (2010), China, and by Midcareer Researcher Program through National Research Foundation (NRF) grant funded by the Ministry of Education, Science and Technology (MEST) (No. 2009-0078682).

REFERENCES

(1) Ben Amara, M.; Vlasse, M.; Le Flem, G.; Hagenmuller, P. Acta Crystallogr, Sect. C 1983, 39, 1483.

(2) Elammari, L.; El Koumiri, M.; Zschokke-Gränacher, I.; Elouadi, B. Ferroelectrics 1994, 158, 19.

(3) Chan, T. S.; Liu, R. S.; Baginskiy, I. Chem. Mater. 2008, 20, 1215.

(4) Lin, C. C.; Xiao, Z. R.; Guo, G. Y.; Chan, T. S.; Liu, R. S. J. Am. Chem. Soc. 2010, 132, 3020.

(5) Zhang, S. Y.; Huang, Y. L.; Seo, H. J. J. Electrochem. Soc. 2010, 157, 261.

(6) Tang, Y. S.; Hu, S. F.; Lin, C. C.; Bagkar, N. C.; Liu, R. S. Appl. Phys. Lett. 2007, 90, 151108.

(7) Im, W. B.; Yoo, H. S.; Vaidyanathan, S.; Kwon, K. H.; Park, H. J.; Kim, Y. I.; Jeon, D. Y. *Mater. Chem. Phys.* **2009**, *115*, 161.

(8) Lin, C. C.; Tang, Y. S.; Hu, S. F.; Liu, R. S. J. Lumin. 2009, 129, 1682.

(9) Wu, Z. C.; Shi, J. X.; Gong, M. L.; Wang, J.; Su, Q. Mater. Chem.

Phys. 2007, 103, 415.
(10) Qin, C. X.; Huang, Y. L.; Shi, L.; Chen, G. Q.; Qiao, X. B.; Seo,
H. J. J. Phys. D: Appl. Phys. 2009, 42, 185105.

(11) Tang, W. J.; Chen, D. H. J. Am. Ceram. Soc. 2009, 92, 1059.

(12) Boutinaud, P.; Parent, C.; Le Flem, G.; Moine, B.; Pedrini, C.

J. Mater. Chem. 1996, 6, 381. (13) Poort, S. H. M.; Blokpoel, W. P.; Blasse, G. Chem. Mater. 1995,

7, 1547. (14) Härknar V.W. Müller Buschhaum H. Alle 7 A Chem 1979.

(14) Hörkner, V. W.; Müller-Buschbaum, H.; Allg, Z. A. Chem. 1979, 451, 40.

(15) Poort, S. H. M.; Reijnhoudt, H. M.; Kuip, H. O. T.; Blasse, G. J. Alloys Compd. **1996**, 241, 75.

(16) Peng, M. Y.; Pei, Z. W.; Hong, G. Y.; Su, Q. J. Mater. Chem. 2003, 13, 1202.

(17) Huang, S.; Von Der Mühll, R.; Ravez, J.; Chaminade, J.; Hagenmuller, P.; Couzi, M. J. Solid State Chem. **1994**, 109, 97.

(18) Liu, B.; Barbier, J. J. Solid State Chem. 1993, 102, 115.

(19) Dougill, M. W. Nature 1957, 180, 292.

(20) Liang, C. S.; Eckert, H.; Gier, T. E.; Stucky, G. D. Chem. Mater. 1993, 5, 597.

(21) Boutinaud, P.; Duloisy, E.; Pedrini, C.; Moine, B.; Parent, C.; Le Flem, G. J. Solid Slate Chem. **1991**, 94, 236.

(22) Wanmaker, W. L.; Spier, H. L. J. Electrochem. Soc. 1962, 109, 109.
(23) Waite, M. S. J. Electrochem. Soc. 1974, 121, 1122.

(24) Wu, Z. C.; Liu, J.; Gong, M. L.; Su, Q. J. Electrochem. Soc. 2009, 156, 153.

(25) Ding, H. Y.; Huang, Y. L.; Shi, L.; Seo, H. J. J. Electrochem. Soc.
 2010, 157, 54.

(26) Elamman, L.; Elouadi, B.; Muller-Vogt, G. Phase Transitions. 1988, 13, 29.

(27) Wu, Z. C.; Liu, J.; Guo, Q. J.; Gong, M. L. Chem. Lett. 2008, 37, 190.

(28) Xia, Q.; Batentschuk, M.; Osvet, A.; Winnacker, A.; Schneider, J. Radiat. Meas. **2010**, *45*, 350.

(29) Inoue, K.; Hirosaki, N.; Xie, R. J.; Takeda, T. J. Phys. Chem. C. 2009, 113, 9392.

(30) Stevels, A. L. N. J. Lumin. 1978, 17, 121.

(31) Bachmann, V.; Ronda, C.; Meijerink, A. *Chem. Mater.* **2009**, *21*, 2077.

(32) Bachmann, V.; Ronda, C.; Oeckler, O.; Schnick, W.; Meijerink, A. *Chem. Mater.* **2009**, *21*, 316.

(33) Srivastava, A. M.; Comanzo, H. A.; Camardello, S.; Chaney, S. B.; Aycibin, M.; Happek, U. J. Lumin. **2009**, *129*, 919.

(34) Poort, S. H. M.; Meijerink, A.; Blasse, G. J. Phys. Chem. Solids 1997, 58, 1451.

(35) Bachmann, V.; Jüstel, T.; Meijerink, A.; Ronda, C.; Schmidt, P. J. J. Lumin. 2006, 121, 441.

(36) Dorenbos, P. J. Phys.: Condens. Matter. 2005, 17, 8103.

(37) Baginskiy, I.; Liu, R. S. J. Electrochem. Soc. 2009, 156, 29.

(38) Yamaga, M.; Masui, Y.; Sakuta, S.; Kodama, N.; Kaminaga, K. *Phys. Rev. B.* **2005**, *71*, 205102.

(39) Vikhnin, V. S.; Liu, G. K.; Beitz, J. V. Phys. Lett. A 2001, 287, 419.

(40) Macfarlane, R. M. Shelby, R. M. In Spectroscopy of Solids

Containing Rare Earth Ions; Kaplyanskii, A. A., Macfarlane, R. M., Eds.; North-Holland: Amsterdam, 1987; p 51.

(41) Meijerink, A.; Dirksen, G. J. J. Lumin. 1995, 63, 189.